

# Twistor geometry

Stephen Huggett\*

September 23, 2022

## Contents

1	Introduction	1
2	Minkowski space	2
3	Twistor space	4
4	Klein correspondence	5
5	Causal structure	6
6	Real points	7
7	Functions on twistor space	8
8	A double fibration	9
9	Sheaf cohomology	10
10	Appendix	12

## 1 Introduction

Our aim is to give a view of the geometrical relationship between Minkowski space, the space-time of special relativity, and projective twistor space, which is  $\mathbb{CP}^3$ . We also give a hint as to how the *Penrose transform* takes functions on twistor space to solutions of field equations in space-time. We refer to [2] for details.

But we warm up by looking at a piece of classical geometry. Given a circle centre  $O$  and radius  $r$  in  $\mathbb{R}^2$ , *inversion* in this circle is the mapping

$$t : \mathbb{R}^2 \setminus \{O\} \rightarrow \mathbb{R}^2 \setminus \{O\}$$

---

\*<http://stephenhuggett.com/index.html>

defined by  $t(A) = A'$ , where  $A'$  lies on the straight line through  $O$  and  $A$ , and on the same side of  $O$  as  $A$ , and  $OA.OA' = r^2$ . This mapping has many lovely properties but is not defined at  $O$ . We can fix this by adding to  $\mathbb{R}^2$  a single point at infinity, as follows.

Place a sphere on a plane, so that the point of contact is the south pole. Let  $P$  be a point on the sphere, other than the north pole  $N$ . Draw a straight line from  $N$  through  $P$ , and suppose that it intersects the plane at  $Q$ . *Stereographic projection* is the mapping

$$\sigma : S^2 \setminus \{N\} \longrightarrow \mathbb{R}^2 \quad (1)$$

$$\sigma(P) \longmapsto Q \quad (2)$$

It is a homeomorphism. The sphere has exactly one extra point,  $N$ , which is “at infinity” in  $\mathbb{R}^2$ .  $S^2$  is called the one-point *compactification* of  $\mathbb{R}^2$ .

Circles and straight lines in  $\mathbb{R}^2$  are mapped by  $\sigma^{-1}$  to circles on the sphere, and, in particular, inversions in circles and reflections in straight lines in  $\mathbb{R}^2$  are well-defined mappings from  $S^2$  to itself. The group of such mappings contains the Euclidean group as a subgroup (because the Euclidean group can be generated by reflections).

In the next section we will see how to define inversions in space-time, and how to compactify space-time so that these inversions are well-defined.

## 2 Minkowski space

A point of space-time is an *event* (which is a point in space and an instant in time). An observer  $S$  in space-time needs four coordinates  $(t, x, y, z)$  to describe the events she sees. Suppose another observer  $S'$ , moving with respect to  $S$ , uses the coordinates  $(t', x', y', z')$ . We choose the  $x'$ ,  $y'$ , and  $z'$  axes to coincide with the  $x$ ,  $y$ , and  $z$  axes when  $t = 0$ , and we also choose  $t' = 0$  then too. Subsequently the origin of  $S'$  moves at constant speed  $v$  along the  $x$  axis of  $S$ , keeping the  $y$  and  $y'$  axes parallel. Note that we also choose  $S$  to be an *inertial* frame of reference, which means that it is not accelerating. Then  $S'$  is also inertial.

Motivated by a desire to unify Newtonian mechanics with Maxwell’s theory of electromagnetism, Einstein adopted the following two physical principles.

- (a) All inertial frames of reference are equivalent for all physical experiments.
- (b) Light has the same velocity  $c$  in all inertial frames of reference.

Einstein discovered that as a direct consequence of these two principles, the transformation between  $S$  and  $S'$  must be given by the *standard Lorentz transformation*

$$t' = \gamma(t - vx/c^2) \quad x' = \gamma(x - vt) \quad y' = y \quad z' = z, \quad (3)$$

where

$$\gamma = \frac{1}{\sqrt{1 - v^2/c^2}}.$$

Minkowski space  $\mathbb{M}$ , the space-time of special relativity, is a four dimensional real vector space. It has a metric

$$dx^a dx^b g_{ab} = (dx^0)^2 - (dx^1)^2 - (dx^2)^2 - (dx^3)^2 \quad (4)$$

where  $x^0 = ct, x^1 = x, x^2 = y, x^3 = z$ . This metric can be positive, zero, or negative, and so at each event

$$x^a = (x^0, x^1, x^2, x^3)$$

in  $\mathbb{M}$  there is a *null cone*, distinguishing these three cases.

For a mathematician, special relativity is the geometry of  $\mathbb{M}$  under the transformations which are isometries of the metric (4). These transformations form the *Lorentz group*  $O(1, 3)$ , which we can enlarge by including translations to form the *Poincaré group*.

**Exercise 1** Show that the transformation (3) is in the Lorentz group.

However, Maxwell's equations for electromagnetism are in fact invariant under an even larger group: the group of *conformal transformations*. This group consists of the Poincaré group together with dilations and inversions.

An inversion in Minkowski space has the formula

$$x^a \mapsto \frac{x^a}{\Delta}, \quad (5)$$

where

$$\Delta = x^a x^b g_{ab}.$$

However, this transformation (5) is not well defined on all of  $\mathbb{M}$ , because it is singular on the null cone  $\Delta = 0$ . To overcome this problem we need to *compactify* Minkowski space.

Consider  $\mathbb{R}^6$  with coordinates  $(T, V, W, X, Y, Z)$  and the metric

$$dT^2 + dV^2 - dW^2 - dX^2 - dY^2 - dZ^2. \quad (6)$$

We use the mapping  $\mathbb{M} \rightarrow \mathbb{R}^6$  given by

$$x^a \mapsto (x^0, \frac{1}{2}(1 - \Delta), -\frac{1}{2}(1 + \Delta), x^1, x^2, x^3) \quad (7)$$

to embed Minkowski space in the intersection of the  $O(2, 4)$  null cone  $N$  of the origin of  $\mathbb{R}^6$ , described by

$$N = \{(T, V, W, X, Y, Z) : T^2 + V^2 - W^2 - X^2 - Y^2 - Z^2 = 0\}, \quad (8)$$

with the hyperplane

$$H = \{(T, V, W, X, Y, Z) : V - W = 1\}. \quad (9)$$

**Exercise 2** Show that the transformation (7) is indeed an embedding of  $\mathbb{M}$  in  $N \cap H$ .

**Exercise 3** Show that on any generator of  $N$  with  $V - W \neq 0$  we can find a point in  $H$  and hence a point of  $\mathbb{M}$ .

**Exercise 4** Show that the transformation (7) is an isometry.

The space of generators of the null cone (8) is a quadric  $Q$  in  $\mathbb{RP}^5$ . It is compact, and is the space which we choose to be compactified Minkowski space  $\mathbb{M}^\#$ . We have added points which correspond to those generators of the null cone which do not intersect  $H$ .

To discover what these extra points look like locally we must go back to the mapping (7). The non-compact space  $\mathbb{M}$  is identified with a subset of the space of generators, isometrically. The extra points in  $\mathbb{M}^\#$  lie on the hyperplane  $V = W$ , and under  $O(2, 4)$  this can be mapped to  $V + W = 0$ , on which  $\Delta = 0$ . So we see that the extra points are a null cone, and we have compactified Minkowski space by adding a null cone at infinity.

Finally, we complexify the space  $\mathbb{M}^\#$  to obtain compactified complexified Minkowski space,  $\mathbb{CM}^\#$ .

### 3 Twistor space

Twistor space,  $\mathbb{T}$ , is a four dimensional complex vector space with complex coordinates

$$Z^\alpha = (Z^0, Z^1, Z^2, Z^3). \quad (10)$$

Projective twistor space,  $\mathbb{PT}$ , is the space of complex lines through the origin in  $\mathbb{T}$ , with homogeneous complex coordinates

$$Z^\alpha = (Z^0 : Z^1 : Z^2 : Z^3). \quad (11)$$

Dual twistor space,  $\mathbb{T}^*$ , is the space of linear functions on  $\mathbb{T}$ :

$$W_\alpha : \mathbb{T} \longrightarrow \mathbb{C} \quad (12)$$

$$Z^\alpha \longmapsto Z^\alpha W_\alpha \quad (13)$$

where we assume summation over the repeated index. Alternatively, if we consider a fixed dual twistor,  $A_\alpha$  say, this defines a hyperplane through the origin in  $\mathbb{T}$  given by

$$\{Z^\alpha : A_\alpha Z^\alpha = 0\}.$$

Dual twistor space is then the space of hyperplanes such as this.

## 4 Klein correspondence

A line in  $\mathbb{PT}$  can be determined by taking the skew-symmetrised outer product of any two planes through it, as follows:

$$L_{\alpha\beta} = X_\alpha Y_\beta - X_\beta Y_\alpha,$$

and then forgetting the overall scale. So the space  $\mathbb{F}_2$  is the space of these  $L_{\alpha\beta}$ , up to an overall scale. In coordinates, this skew-symmetric matrix is

$$L_{\alpha\beta} = \begin{pmatrix} 0 & L_{01} & L_{02} & L_{03} \\ -L_{01} & 0 & L_{12} & L_{13} \\ -L_{02} & -L_{12} & 0 & L_{23} \\ -L_{03} & -L_{13} & -L_{23} & 0 \end{pmatrix}. \quad (14)$$

The space of such matrices up to an overall scale is  $\mathbb{CP}^5$ . However,  $L_{\alpha\beta}$  is not only skew-symmetric, it is also *simple*: it was formed by taking an outer product.

**Lemma 1**  $L_{\alpha\beta}$  (skew) is also simple if and only if

$$L_{\alpha\beta} L_{\gamma\delta} + L_{\alpha\gamma} L_{\delta\beta} + L_{\alpha\delta} L_{\beta\gamma} = 0. \quad (15)$$

**Proof**

Suppose that the skew matrix  $L_{\alpha\beta}$  satisfies (15). Let  $P^\gamma$  and  $Q^\delta$  be arbitrary twistors. Then

$$L_{\alpha\beta} L_{\gamma\delta} P^\gamma Q^\delta + L_{\alpha\gamma} L_{\delta\beta} P^\gamma Q^\delta + L_{\alpha\delta} L_{\beta\gamma} P^\gamma Q^\delta = 0. \quad (16)$$

Now let  $L_{\gamma\delta} P^\gamma Q^\delta = \kappa$ ,  $L_{\alpha\gamma} P^\gamma = X_\alpha$ , and  $L_{\beta\delta} Q^\delta = Y_\beta$ . Then (16) becomes

$$\kappa L_{\alpha\beta} - X_\alpha Y_\beta + X_\beta Y_\alpha = 0, \quad (17)$$

and hence  $L_{\alpha\beta}$  is simple. We leave the other part of the proof as an exercise.  $\square$

**Exercise 5** Finish the proof of Lemma 1.

Therefore the space of skew simple  $L_{\alpha\beta}$  is a quadric in  $\mathbb{CP}^5$ . This is the *Klein correspondence*.

**Theorem 1** The space of lines in  $\mathbb{PT}$  is isomorphic to  $\mathbb{CM}^\#$ .

**Proof**

By making the following convenient choice of coordinates

$$\begin{aligned} T &= \frac{i}{\sqrt{2}}(L_{03} - L_{12}) & V &= L_{23} + \frac{1}{2}L_{01} & W &= L_{23} - \frac{1}{2}L_{01} \\ X &= \frac{i}{\sqrt{2}}(L_{02} - L_{13}) & Y &= \frac{-1}{\sqrt{2}}(L_{02} + L_{13}) & Z &= \frac{-i}{\sqrt{2}}(L_{12} + L_{03}) \end{aligned} \quad (18)$$

equation (15) becomes

$$T^2 + V^2 - W^2 - X^2 - Y^2 - Z^2 = 0$$

which is the null cone of the origin in  $\mathbb{C}^6$ . Recall that we defined  $\text{CM}^\#$  to be the space of generators of this null cone and so our quadric in  $\mathbb{CP}^5$  can be thought of as compactified complexified Minkowski space. Therefore, since this quadric (15) is the space of lines in  $\mathbb{PT}$ , we have established the result.  $\square$

## 5 Causal structure

Let  $x^a$  be the space-time point corresponding to a line  $L \subset \mathbb{PT}$ . For a twistor  $Z^\alpha$  to lie on  $L$  it must satisfy two linear equations. It follows from (7) and (18) that (except when the line is given by  $Z^2 = Z^3 = 0$ ) these two equations can be written

$$\begin{pmatrix} Z^0 \\ Z^1 \end{pmatrix} = \frac{i}{\sqrt{2}} \begin{pmatrix} x^0 + x^3 & x^1 + ix^2 \\ x^1 - ix^2 & x^0 - x^3 \end{pmatrix} \begin{pmatrix} Z^2 \\ Z^3 \end{pmatrix}. \quad (19)$$

**Exercise 6** *Show this.*

Now suppose that the twistor  $Z^\alpha$  also lies on the line corresponding to the space-time point  $y^a$ . Then

$$\begin{pmatrix} Z^0 \\ Z^1 \end{pmatrix} = \frac{i}{\sqrt{2}} \begin{pmatrix} y^0 + y^3 & y^1 + iy^2 \\ y^1 - iy^2 & y^0 - y^3 \end{pmatrix} \begin{pmatrix} Z^2 \\ Z^3 \end{pmatrix}. \quad (20)$$

We will deduce from (19) and (20) that  $w^a = x^a - y^a$  is a null vector. We have

$$\begin{pmatrix} 0 \\ 0 \end{pmatrix} = \frac{i}{\sqrt{2}} \begin{pmatrix} w^0 + w^3 & w^1 + iw^2 \\ w^1 - iw^2 & w^0 - w^3 \end{pmatrix} \begin{pmatrix} Z^2 \\ Z^3 \end{pmatrix}, \quad (21)$$

and therefore

$$0 = \begin{vmatrix} w^0 + w^3 & w^1 + iw^2 \\ w^1 - iw^2 & w^0 - w^3 \end{vmatrix} \quad (22)$$

$$= (w^0 + w^3)(w^0 - w^3) - (w^1 + iw^2)(w^1 - iw^2), \quad (23)$$

so  $w^a$  is null as required.

**Exercise 7** *Show the converse, that if  $x^a - y^a$  is a null vector then the lines in  $\mathbb{PT}$  corresponding to  $x^a$  and  $y^a$  intersect.*

We now have:

$\mathbb{PT}$	$\mathbb{CM}^\#$
complex projective line	point
intersection of lines	null-separation of points

## 6 Real points

The points  $x^a$  and  $y^a$  above are complex: how do we pick out real points in Minkowski space? Any twistor  $Z^\alpha$  lying on the line in  $\mathbb{PT}$  corresponding to the point  $x^a$  satisfies (19), and if  $x^a$  is real then

$$Z^0 \bar{Z}^2 + Z^1 \bar{Z}^3 + \bar{Z}^0 Z^2 + \bar{Z}^1 Z^3 = 0. \quad (24)$$

**Exercise 8** *Show this.*

The Hermitian form

$$\Sigma(Z^\alpha) = Z^0 \bar{Z}^2 + Z^1 \bar{Z}^3 + \bar{Z}^0 Z^2 + \bar{Z}^1 Z^3 \quad (25)$$

divides  $\mathbb{PT}$  into three regions:

$$\begin{aligned} \Sigma(Z^\alpha) > 0 & \quad \mathbb{PT}^+ \\ \Sigma(Z^\alpha) = 0 & \quad \mathbb{PN} \\ \Sigma(Z^\alpha) < 0 & \quad \mathbb{PT}^- \end{aligned}$$

Now go back to (19), and think of  $Z^\alpha$  as a fixed point in  $\mathbb{PN}$ . Any two real solutions  $x^a$  and  $y^a$  of (19) must be null-separated, and so the set of solutions is a (real) null line in Minkowski space. Therefore, points in  $\mathbb{PN}$  correspond to real null lines in  $\mathbb{M}$ . The space of these real null lines is a real 5-dimensional manifold, as can be seen from (25) or from counting them in  $\mathbb{M}$ . Note that not all real 5-dimensional manifolds can be extended to form a complex 3-dimensional manifold: this is a special property of the null lines in  $\mathbb{M}$ .

Also, a line  $L$  lies entirely in the upper half  $\mathbb{PT}^+$  of projective twistor space if and only if it corresponds to a point  $z^a = x^a - iy^a$  in  $\mathbb{CM}^\#$  with  $y^a$  timelike and future-pointing. Fields defined on this part of  $\mathbb{CM}^\#$  have *positive frequency*, an important distinction in quantum field theory which has thus been expressed geometrically in  $\mathbb{PT}$ .

Next, fix a line  $L$  in  $\mathbb{PN}$ , with its corresponding real point  $x^a$  in Minkowski space.  $L$  is a complex projective line  $\mathbb{CP}^1$ , the *Riemann sphere*. Twistors on  $L$  correspond to null lines through  $x^a$ . So intrinsically the line  $L$  in  $\mathbb{PN}$  is the celestial sphere of the space-time point  $x^a$ .

We now have:

PN	$\mathbb{M}^\#$
point	null line
line, intrinsically	celestial sphere of a point

## 7 Functions on twistor space

In this section we consider a simple example of a function on  $\mathbb{P}\mathbb{T}$  and show how to obtain from it a corresponding function on  $\mathbb{C}\mathbb{M}^\#$ . Let

$$f(Z^\alpha) = \frac{1}{(A_\alpha Z^\alpha)(B_\beta Z^\beta)}. \quad (26)$$

This is well defined everywhere in  $\mathbb{P}\mathbb{T}$  except on the two planes  $A_\alpha Z^\alpha = 0$  and  $B_\beta Z^\beta = 0$ . Strictly speaking, it is not a function, because it is homogeneous of degree  $-2$ , not  $0$ . So it should be thought of as a local section of the *sheaf*  $\mathcal{O}(2)$ . We ignore this for the moment, and treat  $f$  as a function.

Choose a point  $x^a \in \mathbb{C}\mathbb{M}^\#$ , and restrict  $f$  to the line in  $\mathbb{P}\mathbb{T}$  corresponding to  $x^a$ . Then from (19)  $A_\alpha Z^\alpha$  becomes

$$A_0 \frac{i}{\sqrt{2}} [(x^0 + x^3)Z^2 + (x^1 + ix^2)Z^3] + A_1 \frac{i}{\sqrt{2}} [(x^1 - ix^2)Z^2 + (x^0 - x^3)Z^3] + A_2 Z^2 + A_3 Z^3,$$

which we abbreviate to  $aZ^2 + bZ^3$ , and similarly for  $B_\beta Z^\beta$ . Now

$$f(Z^\alpha|_{x^a}) = \frac{1}{(aZ^2 + bZ^3)(cZ^2 + dZ^3)}. \quad (27)$$

Note that the complex numbers  $a, b, c$ , and  $d$  depend on  $x^a$ . Next, we perform the following contour integral in the  $\mathbb{C}\mathbb{P}^1$  of the line in  $\mathbb{P}\mathbb{T}$  corresponding to  $x^a$ :

$$\phi(x^a) = \frac{1}{2\pi i} \oint \frac{Z^2 dZ^3 - Z^3 dZ^2}{(aZ^2 + bZ^3)(cZ^2 + dZ^3)} \quad (28)$$

This is done by changing the coordinates

$$\begin{pmatrix} z^0 \\ z^1 \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} Z^2 \\ Z^3 \end{pmatrix} \quad (29)$$

to obtain

$$\phi(x^a) = \frac{1}{2\pi i(ad - bc)} \oint \frac{z^0 dz^1 - z^1 dz^0}{z^0 z^1}. \quad (30)$$



Finally, we put  $z = z^1/z^0$  to map this contour integral down onto  $\mathbb{C}$ , and we obtain

$$\phi(x^a) = \frac{1}{2\pi i(ad - bc)} \oint \frac{dz}{z} \quad (31)$$

$$= \frac{1}{ad - bc}. \quad (32)$$

Suppose the line  $K$  of intersection of the two planes  $Z^\alpha A_\alpha = 0$  and  $Z^\alpha B_\alpha = 0$  corresponds to the point  $k^a$ . Then it can be shown that

$$\frac{1}{ad - bc} = \frac{2}{(A_0 B_1 - A_1 B_0)(x^a - k^a)(x^b - k^b)g_{ab}}. \quad (33)$$

The significance of this result is that  $\phi(x^a)$  is automatically a solution of the *wave equation*

$$\frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} - \frac{\partial^2 \phi}{\partial x^2} - \frac{\partial^2 \phi}{\partial y^2} - \frac{\partial^2 \phi}{\partial z^2} = 0$$

on the complement of the null cone of  $k^a$ .

But the relationship between the twistor function  $f$  and the field  $\phi$  is not one-to-one. There is an equivalence class of twistor functions all of which lead to the same  $\phi$ . We will see that this class is described by Čech cohomology, as follows.

**Theorem 2**

$$\check{H}^1(\mathbb{P}\mathbb{T} \setminus K; \mathcal{O}(-2)) \cong \{ \text{analytic solutions of the wave equation on the complement of the null cone of } k^a \} \quad (34)$$

In fact, all analytic solutions of all the *zero rest mass free field* equations (which include Maxwell's equations) can be constructed in a similar way, by starting with functions  $f(Z^\alpha)$  homogeneous of degree  $n$ , for various integers  $n$ .

## 8 A double fibration

Here we describe the geometrical framework behind the calculation of  $\phi(x^a)$  from  $f(Z^\alpha)$ . Define

$$\begin{aligned} \mathbb{F}_1 &= \{S_1 : S_1 \text{ is a one-dimensional subspace of } \mathbb{T}\} \\ \mathbb{F}_2 &= \{S_2 : S_2 \text{ is a two-dimensional subspace of } \mathbb{T}\} \\ \mathbb{F}_{1,2} &= \{(S_1, S_2) : S_1 \text{ and } S_2 \text{ as above with } S_1 \text{ a subspace of } S_2\}. \end{aligned}$$

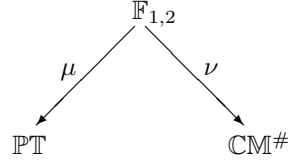
Note that  $F_1 = \mathbb{P}\mathbb{T}$  and  $F_2 = \mathbb{C}\mathbb{M}^\#$ . Now define the “forgetful” maps  $\mu : \mathbb{F}_{1,2} \longrightarrow \mathbb{F}_1$  given by

$$\mu(S_1, S_2) = S_1 \quad (35)$$

and  $\nu : \mathbb{F}_{1,2} \longrightarrow \mathbb{F}_2$  given by

$$\nu(S_1, S_2) = S_2, \quad (36)$$

and we then have the double fibration



**Exercise 9** For  $Z^\alpha$  a point in  $\mathbb{P}\mathbb{T}$ , describe  $\mu^{-1}(Z^\alpha)$ , which is called the fibre of  $\mu$  above  $Z^\alpha$ .

**Exercise 10** For  $x^a$  a point in  $\mathbb{C}\mathbb{M}^\#$ , describe  $\nu^{-1}(x^a)$ , which is called the fibre of  $\nu$  above  $x^a$ .

Now the procedure in the previous section can be expressed as follows. Given a function  $f(Z^\alpha)$  on a subset of  $\mathbb{P}\mathbb{T}$ , we define its *pullback*

$$\mu^*(f)(Z^\alpha, x^a) = f(Z^\alpha), \quad (37)$$

which is a function on a subset of  $\mathbb{F}_{1,2}$ . Then we integrate  $\mu^*(f)$  along the fibres of  $\nu$ , to obtain the function  $\phi(x^a)$  on a subset of  $\mathbb{C}\mathbb{M}^\#$ . What are the conditions on these subsets? In general, we start by choosing

$$X \subset \mathbb{C}\mathbb{M}^\#$$

and then select the corresponding spaces

$$Y = \nu^{-1}(X) \subset \mathbb{F}$$

and

$$T = \mu(Y) \subset \mathbb{P}\mathbb{T}.$$

The procedure for obtaining  $\phi$  from  $f$  will work when the fibres of  $\mu|_Y$  are connected and simply connected.

## 9 Sheaf cohomology

This is an *extremely* brief description of sheaf cohomology. See [1] for a proper account.

Given a topological space  $X$ , a *sheaf* over  $X$  is another topological space  $\mathcal{S}$  with a mapping  $\pi : \mathcal{S} \rightarrow X$  satisfying

1.  $\pi$  is a local homeomorphism,
2. the *stalks*  $\pi^{-1}(x)$  are abelian groups, and
3. the group operations are continuous.

If  $U$  is open in  $X$  we denote by  $\mathcal{S}(U)$  the abelian group of *sections* of  $\mathcal{S}$  over  $U$ .

As an example, take holomorphic functions on a complex manifold. Consider a point  $z$  in the manifold and a function element  $(f, D)$  such that the domain  $D$  contains  $z$ . We define the *germ*  $[f, z]$  of  $f$  at  $z$  to be *all* function elements  $(f_i, D_i)$  such that  $z \in D_i$  and there is a neighbourhood of  $z$  on which  $f_i = f$ . Then the set of germs is the sheaf  $\mathcal{O}$  on the manifold.

Now we turn to Čech cohomology, still with our sheaf  $\mathcal{S}$ . Choose an open cover  $\{U_i\}$  of  $X$ , and simplify the notation for intersections by setting

$$U_{ij} = U_i \cap U_j. \quad (38)$$

A 0-cochain is a collection

$$\{f_i \in \mathcal{S}(U_i)\}, \quad (39)$$

a 1-cochain is a collection

$$\{f_{ij} \in \mathcal{S}(U_{ij}) : f_{ij} = -f_{ji}\}, \quad (40)$$

a 2-cochain is a collection

$$\{f_{ijk} \in \mathcal{S}(U_{ijk}) : f_{ijk} \text{ is skew symmetric}\}, \quad (41)$$

and similarly for a  $p$ -cochain. The set of all  $p$ -cochains forms an abelian group which we denote by  $C^p(\{U_i\}; \mathcal{S})$ .

We have a *coboundary map*  $\delta_p : C^p(\{U_i\}; \mathcal{S}) \rightarrow C^{p+1}(\{U_i\}; \mathcal{S})$ , a group homomorphism defined as follows:

$$\delta_0(\{f_i\}) = \{\rho_j f_i - \rho_i f_j\}, \quad (42)$$

$$\delta_1(\{f_{ij}\}) = \{\rho_k f_{ij} + \rho_j f_{ki} + \rho_i f_{jk}\}. \quad (43)$$

Here,  $\rho_i$  means restrict to the intersection with  $U_i$ .  $\delta_p$  is defined similarly, always in such a way that the result is skew-symmetric. It is easy to check that  $\delta_1 \circ \delta_0 = 0$ , and in fact the skew symmetry implies that  $\delta_{p+1} \circ \delta_p = 0$  for all  $p$ .

This gives us a sequence of abelian groups

$$C^0(\{U_i\}; \mathcal{S}) \xrightarrow{\delta_0} C^1(\{U_i\}; \mathcal{S}) \xrightarrow{\delta_1} C^2(\{U_i\}; \mathcal{S}) \xrightarrow{\delta_2} \dots \quad (44)$$

in which the image of  $\delta_{p-1}$  is a normal subgroup of the kernel of  $\delta_p$ . We define

$$\check{H}^p(\{U_i\}; \mathcal{S}) = \frac{\ker \delta_p}{\text{im } \delta_{p-1}}. \quad (45)$$

Strictly speaking we should now take successive refinements of the cover  $\{U_i\}$  of  $X$  and then the direct limit, in order to remove the dependence on any particular cover and obtain  $\check{H}^p(X; \mathcal{S})$ , the Čech cohomology of  $X$  with coefficients in the sheaf  $\mathcal{S}$ . But in practice it is usually possible to find a specific cover such that

$$\check{H}^p(\{U_i\}; \mathcal{S}) = \check{H}^p(X; \mathcal{S}). \quad (46)$$

Looking back at the contour integral we carried out, we can now see that the equivalence class of functions all leading to the same field  $\phi$  is simply an element of  $\check{H}^1(\{U_1, U_2\}; \mathcal{O}(-2))$ , where  $U_1$  is the complement of the plane  $A_\alpha Z^\alpha = 0$  and  $U_2$  is the complement of the plane  $B_\alpha Z^\alpha = 0$ .

More generally, it can be shown that

**Theorem 3**

$$\check{H}^1(\mathbb{P}\mathbb{T}^+; \mathcal{O}(-2)) \cong \{ \text{positive-frequency analytic solutions} \\ \text{of the wave equation} \} \quad (47)$$

## 10 Appendix

**Exercise 1.** Using (3), calculate

$$(ct')^2 - (x')^2 - (y')^2 - (z')^2,$$

and simplify.

**Exercise 2.** Firstly, note that if

$$(x^0, \frac{1}{2}(1 - \Delta), -\frac{1}{2}(1 + \Delta), x^1, x^2, x^3) = (y^0, \frac{1}{2}(1 - \Delta), -\frac{1}{2}(1 + \Delta), y^1, y^2, y^3)$$

then

$$x^a = y^a.$$

So the mapping is injective. Clearly, it is also continuous.

**Exercise 3.** A generator of  $N$  is the set of points

$$\{(\lambda T, \lambda V, \lambda W, \lambda X, \lambda Y, \lambda Z) : \lambda \neq 0 \in \mathbb{C}\}.$$

If  $V - W \neq 0$  we can choose  $\lambda = (V - W)^{-1}$ , and then this point on the generator will have values of  $V$  and  $W$  satisfying  $V - W = 1$ .

**Exercise 4.** Substitute  $V - W = 1$  into (6).

**Exercise 5.** Suppose that  $L_{\alpha\beta}$  is simple. Then

$$L_{\alpha\beta} = X_\alpha Y_\beta - X_\beta Y_\alpha.$$

Now insert this expression into the left hand side of (15) and simplify.

**Exercise 6.** Equation (19) can be split into the two equations

$$Z^\alpha X_\alpha = 0 \quad \text{and} \quad Z^\alpha Y_\alpha = 0,$$

where

$$\begin{aligned} X_\alpha &= \left[ -1, 0, \frac{i}{\sqrt{2}}(x^0 + x^3), \frac{i}{\sqrt{2}}(x^1 + ix^2) \right] \\ Y_\alpha &= \left[ 0, -1, \frac{i}{\sqrt{2}}(x^1 - ix^2), \frac{i}{\sqrt{2}}(x^0 - x^3) \right]. \end{aligned}$$

Now set

$$L_{\alpha\beta} = X_\alpha Y_\beta - X_\beta Y_\alpha,$$

and then show that

$$\frac{i}{\sqrt{2}}(L_{03} - L_{12}) = x^0, \quad \frac{i}{\sqrt{2}}(L_{02} - L_{13}) = x^1,$$

and so on, as required for (7) and (18).

**Exercise 7.** Suppose that  $w^a = x^a - y^a$  is a null vector. Then the rows in

$$\begin{pmatrix} w^0 + w^3 & w^1 + iw^2 \\ w^1 - iw^2 & w^0 - w^3 \end{pmatrix}$$

are linearly dependent, so we can find  $Z^2$  and  $Z^3$  satisfying (21). Now calculate  $Z^0$  and  $Z^1$  from (19) or (20). Then  $Z^\alpha$  lies on both the line corresponding to  $x^a$  and the line corresponding to  $y^a$ .

**Exercise 8.**

$$\begin{aligned} Z^0 \overline{Z^2} &= \frac{i}{\sqrt{2}} \left[ (x^0 + x^3) Z^2 \overline{Z^2} + (x^1 + ix^2) Z^3 \overline{Z^2} \right] \\ Z^1 \overline{Z^3} &= \frac{i}{\sqrt{2}} \left[ (x^1 - ix^2) Z^2 \overline{Z^3} + (x^0 - x^3) Z^3 \overline{Z^3} \right] \\ \overline{Z^0} Z^2 &= \frac{-i}{\sqrt{2}} \left[ (x^0 + x^3) Z^2 \overline{Z^2} + (x^1 - ix^2) Z^2 \overline{Z^3} \right] \\ \overline{Z^1} Z^3 &= \frac{-i}{\sqrt{2}} \left[ (x^1 + ix^2) Z^3 \overline{Z^2} + (x^0 - x^3) Z^3 \overline{Z^3} \right] \end{aligned}$$

and hence the result.

**Exercise 9.**  $\mu^{-1}(Z^\alpha)$  is the space of all pairs  $\{Z^\alpha, L\}$  where  $L$  is a line in  $\mathbb{PT}$  which goes through the point  $Z^\alpha$ . So it is the space of all pairs  $\{Z^\alpha, x^a\}$ , where  $x^a$  satisfies (19) for the given  $Z^\alpha$ .

**Exercise 10.**  $\nu^{-1}(x^a)$  is the space of all pairs  $\{Z^\alpha, x^a\}$  where  $Z^\alpha$  lies on the line in  $\mathbb{PT}$  corresponding to  $x^a$ . So it is the space of all pairs  $\{Z^\alpha, x^a\}$ , where  $Z^\alpha$  satisfies (19) for the given  $x^a$ .

## References

- [1] P Griffiths and J Harris 1978 *Principles of Algebraic Geometry*, Wiley.
- [2] S A Huggett and K P Tod 1994 *An introduction to twistor theory*, second edition, Cambridge University Press.